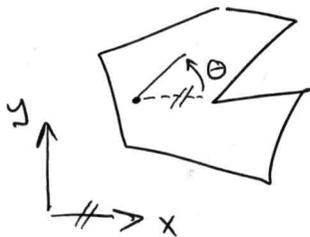


4.2: Defining the C-space4.2.1: C-space of rigid bodies in the plane: $SE(2)$

Goal: characterize all ^{rigid body} displacements

Valid config variables: If the variables are known, then every point of the rigid body is known.

One choice is (x, y, θ) .



(x, y) = position of ref. pt.

θ = angle of ref. line

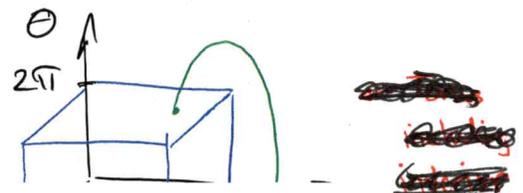
elements of $SE(2) \Rightarrow \begin{bmatrix} c_\theta & -s_\theta & x \\ s_\theta & c_\theta & y \\ 0 & 0 & 1 \end{bmatrix}$

C-space for (x, y) is $M_1 = \mathbb{R}^2$, where M_1 is a manifold

C-space for θ is $M_2 = S^1$ (the circle), also a manifold

$C = M_1 \times M_2 = \mathbb{R}^2 \times S^1$, product of manifolds is a manifold

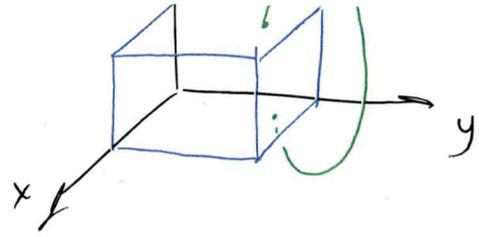
If $M_1 \subset \mathbb{R}^2$, then \rightarrow



if x, y, θ are bounded, then

$$C = I' \times I' \times S'$$

C is a solid torus.

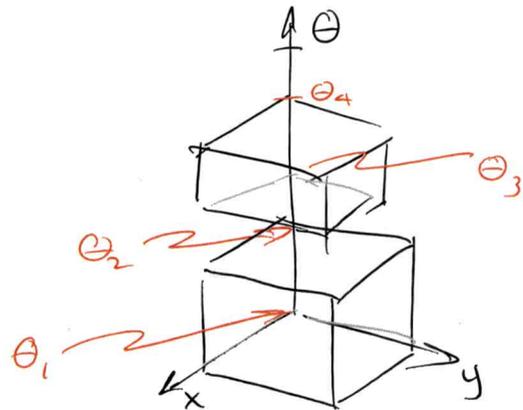


If θ is also bounded, then $C = I' \times I' \times I' = I^3$
 which is a solid sphere or cube

If $\theta \in \{(\overbrace{\theta_1, \theta_2}^{I'}, \overbrace{\theta_3, \theta_4}^{I'})\}$, then

$$C = I^3 \sqcup I^3$$

disjoint union

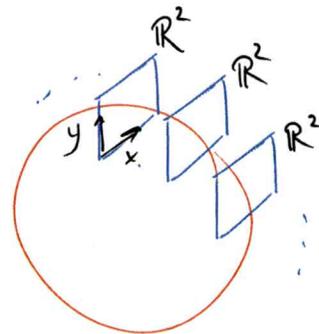


If x, y, θ are unbounded,

then we have a

"fibration" of a circle

where the fibers are \mathbb{R}^2



Matrix Groups :

Main interest is in $SO(n)$ = all $n \times n$ rotation matrices

General Linear Group: $GL(n)$

The space of all $n \times n$ real-valued nonsingular matrices

$$\begin{aligned} A \in GL(n) & \text{ iff } A^{-1} \text{ exists} \\ & \equiv \text{ iff } \det(A) \neq 0 \end{aligned}$$

$$\text{Dim}(GL(n)) = n^2 \quad \text{Think of as } \mathbb{R}^{n^2} \setminus \{\det(A)=0\}$$

An n^2 -dim set w/ a 1-dimensional set removed.

Orthogonal Group: $O(n)$

$$\begin{aligned} A \in O(n) & \text{ iff } A \in GL(n) \text{ and } A^T = A^{-1} \\ & \text{ie., } A^T A = I \end{aligned}$$

\therefore all pairs of columns are orthogonal
and all columns have unit length.

$$\text{Dim}(O(n)) = n^2 - \binom{n}{2} - n = n^2 - \frac{(n-1)n}{2} - n$$

$$= \frac{n(n-1)}{2}$$

of orthogonality relationships
of normality relationships

$\text{Dim}(O(2)) = 1 = \# \text{ dof of orientation in plane}$

$\text{Dim}(O(3)) = 3 = \# \text{ dof " " " } \mathbb{R}^3$

If $A \in O(n)$, then $\text{Det}(A) = 1$ or -1 .

half have 1, half have -1.

Special Orthogonal Group: $SO(n)$

$A \in SO(n)$ iff $A \in O(n)$ and $\text{det}(A) = +1$.

$\text{Dim}(SO(2)) = 1$

$\text{Dim}(SO(3)) = 3$

If space was 4D ... $\text{Dim}(SO(4)) = 6$

constraint apparently
wipes out half of
the space of $SO(n)$,
so dimension is
unchanged.

Example: Matrix groups.

$\left\{ \begin{bmatrix} a & c \\ b & d \end{bmatrix} \mid (a, b, c, d) \in \mathbb{R}^4 \right\}$ is homeomorphic
to \mathbb{R}^4

$GL(2) = \left\{ \begin{bmatrix} a & c \\ b & d \end{bmatrix} \mid (a, b, c, d) \in \mathbb{R}^4, ad - bc \neq 0 \right\}$

$\text{Dim}(GL(2)) = 4$, but is not homeo. to \mathbb{R}^4

because we've removed a one-dimension set
(which has no ~~area~~!)
Volume

Next enforce $A^T A = I$

$$O(2) \supset GL(2) \Rightarrow \begin{bmatrix} a & c \\ b & d \end{bmatrix} \begin{bmatrix} a & b \\ c & d \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$$

$$3 \text{ eqs. } \left\{ \begin{array}{l} a^2 + c^2 = 1 \\ ab + cd = 0 \\ ba + dc = 0 \\ b^2 + d^2 = 1 \end{array} \right\} \text{ same} \Rightarrow \text{Dim}(O(2)) = 1$$

What is the structure of $O(2)$?

$$\text{Let } \begin{bmatrix} a \\ b \end{bmatrix} = w, \|w\| = 1; \begin{bmatrix} c \\ d \end{bmatrix} = v, \|v\| = 1$$

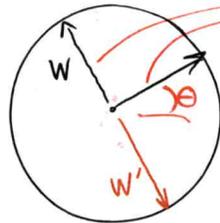
$$w \perp v$$

Clearly v can move on a circle, S^1 , possible (a,b) pairs

Let S^1 be the base manifold.

For each pt. on S^1 there are 2 distinct w 's

so the space $O(2)$ is $S^1 \sqcup S^1$, i.e.



(This comment is better for $SO(2)$)
these can be thought of as moving as a rigid unit to trace out a circle.

Only $\theta \in [0, 2\pi)$ needed to parametrize

disjoint union
union operation that retains index to set of original membership

two circles that never touch ($w \neq w'$ for any v).

$SO(2)$ requires $ad-bc=1$ ← This eq. selects one of the two circles of $O(2)$.

∴ $SO(2)$ is one circle, $SO(2)$ is homeo. to S^1

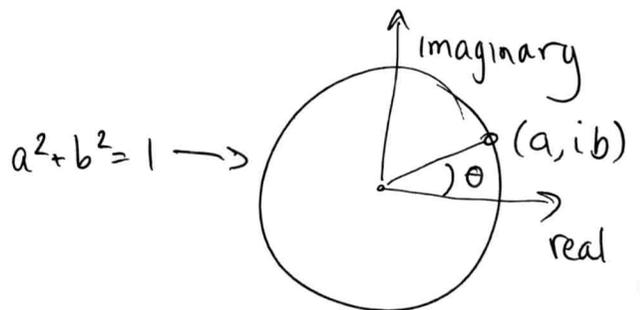
Interpret $SO(2)$ as a circle in the complex plane

$$a+ib \equiv re^{i\theta}$$

$$i^2 = -1$$

$$r = \sqrt{a^2+b^2}$$

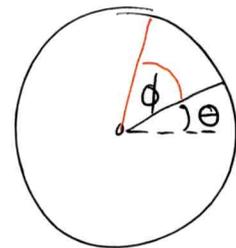
$$\theta = \text{atan2}(b, a)$$



(of unit length)

Note that multiplying a complex # by another \wedge has effect of rotating the first.

$$re^{i\theta} \cdot \cancel{re^{i\phi}} = re^{i(\theta+\phi)}$$



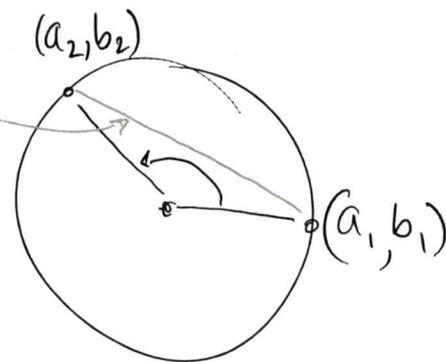
Unit quaternions work similarly in $SO(3)$!

Skip to page 2/7/8

Linear interpolation of angles?

$$\begin{bmatrix} a \\ b \end{bmatrix} = \alpha \begin{bmatrix} a_1 \\ b_1 \end{bmatrix} + (1-\alpha) \begin{bmatrix} a_2 \\ b_2 \end{bmatrix}$$

Not valid rotation.



Use SLERP (Spherical Linear Interpolation) to stay on circle in $SO(2)$ & 3-sphere in $SO(3)$.

$$v = \text{slerp}(\alpha, \hat{v}_1, \hat{v}_2) = \frac{\sin((1-\alpha)\phi)}{\sin(\phi)} \hat{v}_1 + \frac{\sin(\alpha\phi)}{\sin(\phi)} \hat{v}_2$$

where $\phi = \cos(\hat{v}_1 \cdot \hat{v}_2)$

This works in $SO(2)$ & $SO(3)$.

$$\alpha = 0 \Rightarrow v = \hat{v}_1$$

$$\alpha = 1 \Rightarrow v = \hat{v}_2$$

There is no analogous simple formula for Euler angles (or other 3-param. representations of $SO(3)$).

Special Euclidean Group: $SE(n)$

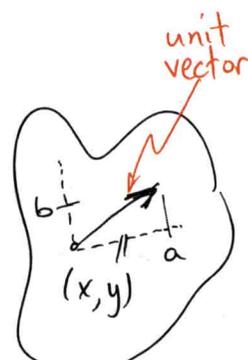
$$SE(2) = \mathbb{R}^2 \times SO(2)$$

one parametrization

$$w/ a^2 + b^2 = 1$$

$$\begin{bmatrix} a & -b & x \\ b & a & y \\ 0 & 0 & 1 \end{bmatrix}$$

2 parameters



When orientation is updated incrementally, $a^2 + b^2 \neq 1$.

Make sure to normalize after rotation change.

$$w/ a^2 + b^2 = 1$$

$$\begin{bmatrix} 0 & 0 & 1 \end{bmatrix}$$



another (θ, x, y)



1-param. rep. of $SO(2)$

$SE(3) = ?$

parameterization: $(x, y, z, \text{roll}, \text{pitch}, \text{yaw})$
 $\underbrace{\mathbb{R}^3}_{\text{6 params}} \quad \underbrace{S^3 = S' \times S' \times S' ?}_{\text{6 params}}$

another param. : (x, y, z, Q)
 $\underbrace{\text{7 params}}_{\text{LaValle notation. unit quat } \triangleq h}$
 where $Q = \text{unit quaternion}$

$$Q = (a, b, c, d) \text{ with } a^2 + b^2 + c^2 + d^2 = 1$$

$SE(3) = \mathbb{R}^3 \times SO(3)$ NOT! $\mathbb{R}^3 \times S^1 \times S^1 \times S^1$
 $S^1 = T^1$ (the 1-torus)

Problem with Euler Angle (all 3-angle) parameterizations!

singularities: $(\alpha, \beta, \gamma) = (0, 0, 0) \Rightarrow I_{3 \times 3} = \text{no rotation}$

However, $(\alpha, \beta, \gamma) \neq (0, 0, 0)$ can also correspond to no rotation

This is a problem in motion planning, since we need to use the inverse orientation mapping a lot.

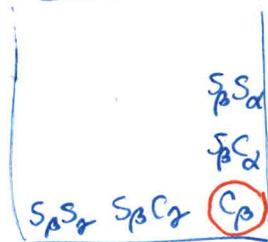
~~For any singularity~~, a continuum of (α, β, γ) exists.
 For for any ~~reachable~~ orientation singular

$$Rot_x(\alpha) Rot_y(\beta) Rot_z(\gamma)$$

e.g. for X-Y-X Euler angles, we have these problems whenever β is near 0 or π .

Also, when near a singular config, numerical solution loses accuracy.

Every 3 angle rep has singularities Euler, RPY, ... in some order



$(\alpha, \beta, \gamma) \rightarrow R$ (rotation matrix) is easy, but

$R \rightarrow (\alpha, \beta, \gamma)$ can be difficult

← we did this for spherical wrist IK solution of Puma

$C_\beta = \text{const}$ becomes singular when $C_\beta = \pm 1$.

THE GROUP OF UNIT QUATERNIONS, H.

Unit Quaternions (a.k.a. Euler Parameters)

$Q \rightarrow R$ easy

$R \rightarrow Q$ easy

Lavelle uses h.

Q can be thought of as a complex #

$$Q = a + bi + cj + dk$$

$$i^2 = j^2 = k^2 = -1$$

$$ij = k, jk = i, ki = j$$

Complex conjugate: $Q^* = a - bi - cj - dk$

$$Q_1 \cdot Q_2 = Q_3$$

$$= (a_1 + b_1 i + c_1 j + d_1 k) \cdot (a_2 + b_2 i + c_2 j + d_2 k)$$

$$a_3 = a_1 a_2 - b_1 b_2 - c_1 c_2 - d_1 d_2$$

$$b_3 = a_1 b_2 + a_2 b_1 + c_1 d_2 - c_2 d_1$$

$$c_3 = \text{similar}$$

$$d_3 = \text{similar}$$

} eq 4.19, pg 150 of
Lavelle

Note that (unit quat) · (unit quat) = (unit quat)

There is no equivalent way to multiply Euler angles (or elements of other 3-param reps).

(Unless you view $R(\alpha, \beta, \gamma) R(\alpha_2, \beta_2, \gamma_2)$ as equivalent!)

Alternative view

$$Q = (a, \mathbf{n})$$

vector part

scalar part

$$Q_1 \cdot Q_2 = (a_1, \mathbf{n}_1) \cdot (a_2, \mathbf{n}_2)$$

$$Q_1 \cdot Q_2 = ((a_1 a_2 - \mathbf{n}_1 \cdot \mathbf{n}_2), (a_1 \mathbf{n}_2 + a_2 \mathbf{n}_1 + \mathbf{n}_1 \times \mathbf{n}_2))$$

usual dot prod.

Think of multiplying these in all possible pairs. Some give scalars, some give vectors.

Mapping Q to R

$$\begin{bmatrix} 2(a^2 + b^2) - 1 & 2(bc - ad) & 2(bd + ac) \end{bmatrix}$$

$$R(Q) = \begin{bmatrix} 2(a^2+b^2)-1 & 2(bc-ad) & 2(bd+ac) \\ 2(bc+ad) & 2(a^2+c^2)-1 & 2(cd-ab) \\ 2(bd-ac) & 2(cd+ab) & 2(a^2+d^2)-1 \end{bmatrix}$$

Mapping $R \rightarrow Q$

$$\text{Let } R = \begin{bmatrix} r_{11} & r_{12} & r_{13} \\ r_{21} & r_{22} & r_{23} \\ r_{31} & r_{32} & r_{33} \end{bmatrix} \text{ and}$$

replacing 1 in $R(Q)$ with $a^2+b^2+c^2+d^2$ yields:

$$a^2 = \frac{1}{4}(1 + r_{11} + r_{22} + r_{33}) \quad (1)$$

$$b^2 = \frac{1}{4}(1 + r_{11} - r_{22} - r_{33}) \quad (2)$$

$$c^2 = \frac{1}{4}(1 - r_{11} + r_{22} - r_{33}) \quad (3)$$

$$d^2 = \frac{1}{4}(1 - r_{11} - r_{22} + r_{33}) \quad (4)$$

These equations imply 16 solutions, but only 2 exist.

Off-diagonal terms give:

$$(5) \quad ab = \frac{1}{4}(r_{32} - r_{23}) \quad bc = \frac{1}{4}(r_{12} + r_{21}) \quad (8)$$

$$(6) \quad ac = \frac{1}{4}(r_{13} - r_{31}) \quad bd = \frac{1}{4}(r_{13} + r_{31}) \quad (9)$$

$$(7) \quad ad = \frac{1}{4}(r_{21} - r_{12}) \quad cd = \frac{1}{4}(r_{23} + r_{32}) \quad (10)$$

Soln approach (to minimize numerical error):

- ① Use equation (1), (2), (3), or (4) with largest value to get 2 values of $a, b, c,$ or d .
- ② Use 3 eqs from (5-10) to solve for other 3 values
- ③ Result is 2 solutions satisfying $Q_1 = -Q_2$

Alg.

If $\text{trace}(R) > 0$

solve eq (1) for $a = \frac{1}{2} \sqrt{\text{tr}(R) + 1}$

solve for b, c, d using equations (5-7)

$$\text{eq(5) yields } \Rightarrow b = \frac{r_{32} - r_{23}}{4a}$$

eq(6) \Rightarrow similar

eq(7) \Rightarrow similar

else if $r_{11} = \max(r_{11}, r_{22}, r_{33})$

solve for b using eq(2)

solve for a, c, d using eqs (5, 8, 9)

else if $r_{22} = \max(r_{11}, r_{22}, r_{33})$

solve for c using eq(3)

solve for a, b, d using eqs (6, 8, 10)

else

else

solve for d using eq (4)

solve for a, b, c using eqs (7, 9, 10)

end

finally $Q_1 = (a, b, c, d) \neq Q_2 = -Q_1$

are the two solutions.

Note: The group of unit quaternions covers $SO(3)$

twice w/ no point identified between the covers,

i.e. Q is never equal to $-Q$, since $Q=0 \notin H$

Step $\frac{\sin((1-a)\phi)}{\sin(\alpha)} Q_1 + \frac{\sin(\alpha\phi)}{\sin(\alpha)} Q_2 = Q$ (rotation about fixed axis)

space of
unit Quats

Rotating a point w/ Q

$$R(Q) \begin{bmatrix} x \\ y \\ z \end{bmatrix} \equiv (h + p + h^*)_{\text{vector part}}$$

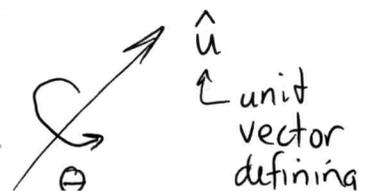
$$\text{where } p = \begin{bmatrix} x \\ y \\ z \\ 0 \end{bmatrix}$$

$$h^* = a - bi - cj - dk$$

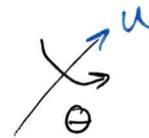
Do not confuse axis/angle with unit quaternion!

axis/angle $\rightarrow Q$

$$Q = \left(\cos\left(\frac{\theta}{2}\right), \hat{u} \sin\left(\frac{\theta}{2}\right) \right)$$

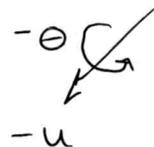


$$Q = (\cos(\frac{\theta}{2}), \hat{u} \sin(\frac{\theta}{2}))$$



vector
defining
axis of
rotation.

Easy to see that $Q = -Q$



$Q \rightarrow$ axis/angle

$$\theta = 2 \cdot \text{atan2}(\|N\|, a)$$

$$u = N / \|N\|$$

A unit quaternion is NOT the same as an axis/angle

$$R(\theta, u) = \begin{bmatrix} u_x^2 + c_\theta(1-u_x^2) & u_x u_y(1-c_\theta) + u_z s_\theta & u_x u_z(1-c_\theta) - u_y s_\theta \\ u_x u_y(1-c_\theta) - u_z s_\theta & u_y^2 + c_\theta(1-u_y^2) & u_y u_z(1-c_\theta) + u_x s_\theta \\ u_x u_z(1-c_\theta) + u_y s_\theta & u_y u_z(1-c_\theta) - u_x s_\theta & u_z^2 + c_\theta(1-u_z^2) \end{bmatrix}$$

Note that u is the eigenvector of R corresponding to the eigenvalue 1

Unit quaternions have a natural metric defining distance between Q_1 & Q_2 . It is the length of the arc of great circle they define on the 3-sphere S^3 .

$$\alpha = \cos^{-1}(Q_1 \cdot Q_2)$$

↑ usual dot product of vectors

A metric allows us to construct algorithms to perform uniform sampling of $SO(3)$.

Velocity Kinematics of Unit Quaternions

$$\begin{bmatrix} \dot{a} \\ \dot{b} \\ \dot{c} \\ \dot{d} \end{bmatrix} = \frac{1}{2} \underbrace{\begin{bmatrix} -b & -c & -d \\ a & -d & c \\ d & a & -b \\ -c & b & a \end{bmatrix}}_B \begin{bmatrix} \omega_x \\ \omega_y \\ \omega_z \end{bmatrix} = \frac{1}{2} B \omega$$

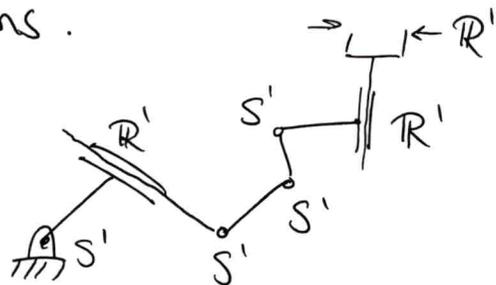
where ω is angular velocity of body expressed in the body-fixed frame.

$$\dot{Q} = B(Q) \omega$$

After incremental rotation ($Q + \dot{Q} \Delta t = Q_{new}$), $\|Q_{new}\| \neq 1$, so normalize

Finally apply to robot systems.

1 dof joints
w/o limits \Rightarrow



\mathbb{R}^1 or S^1

1 dof joints w/ limits
 $\Rightarrow I^1$

$$C = S^1 \times \mathbb{R}^1 \times (S^1)^3 \times \mathbb{R}^2$$

$$C = S^4 \times \mathbb{R}^3$$

with limits

$$C = I^7$$

For every free rigid body, cross C with $\mathbb{R}^3 \times SO(3)$.

For every spherical joint, cross C with $SO(3)$.

Remember: Product of manifolds is a manifold, so C -spaces \neq
State spaces robots are manifolds (unless they have
closed kinematic chains)